MESA+

INSTITUTE FOR NANOTECHNOLOGY

UNIVERSITY OF TWENTE.

Josephson supercurrent through a topological insulator surface state

Nb/Bi_2Te_3/Nb junctions

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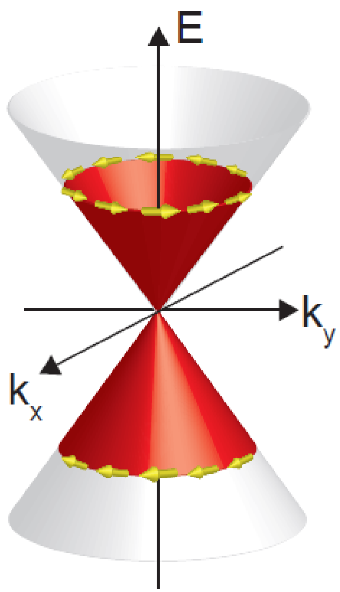
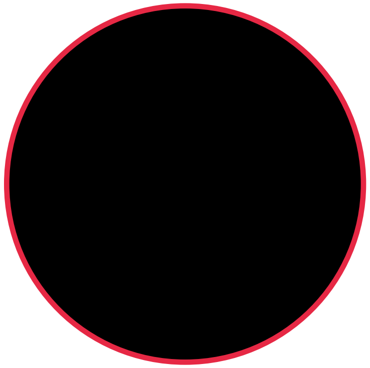
Motivation

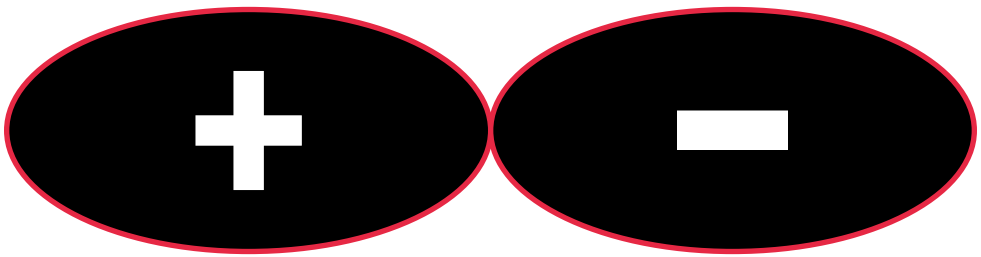


TI surface states



S-wave
Superconductor

Motivation



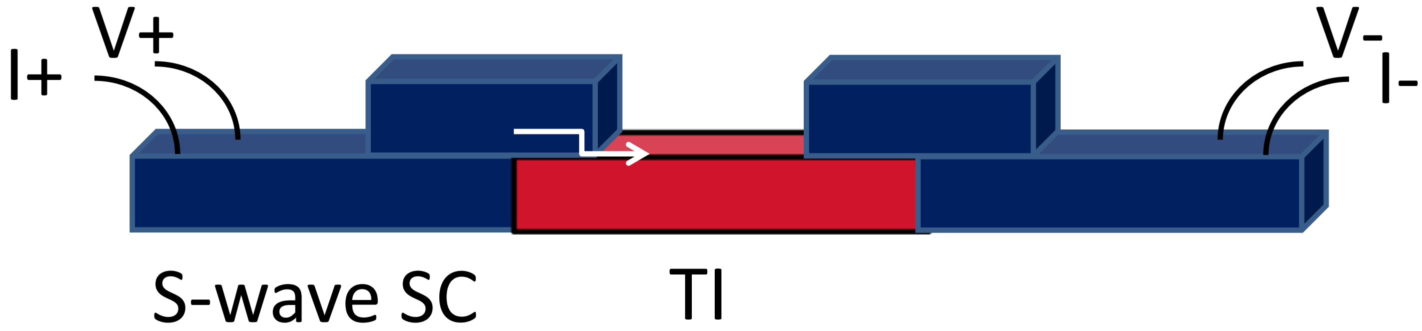
P-wave Superconductor

Natural place to look for Majorana fermions (single zero-energy modes)

Motivation

Essential: supercurrent must couple to surface states

Characterize junction

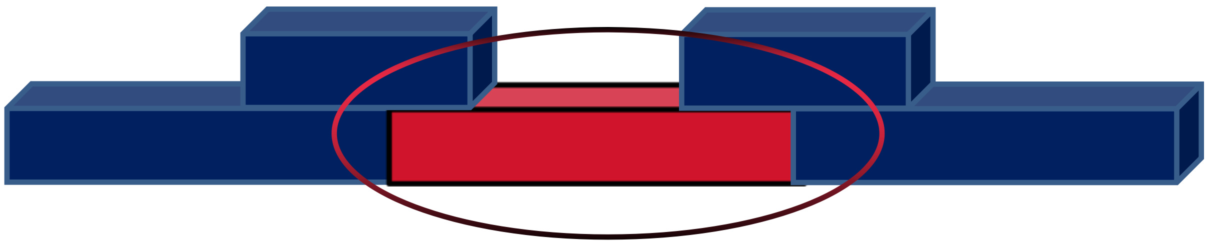


Content

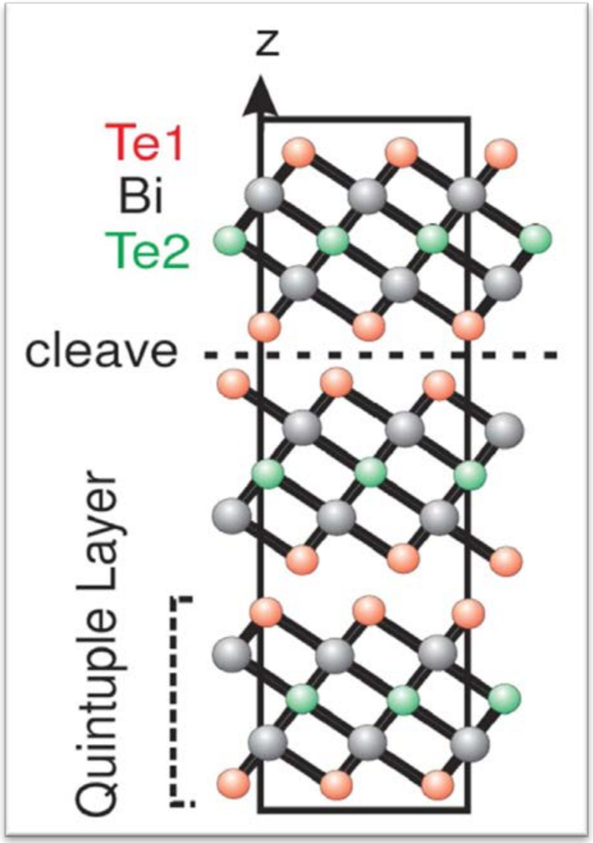
Part 1	Bi_2Te_3
Part 2	S/TI/S junctions
Part 3	Josephson supercurrent through the surface states

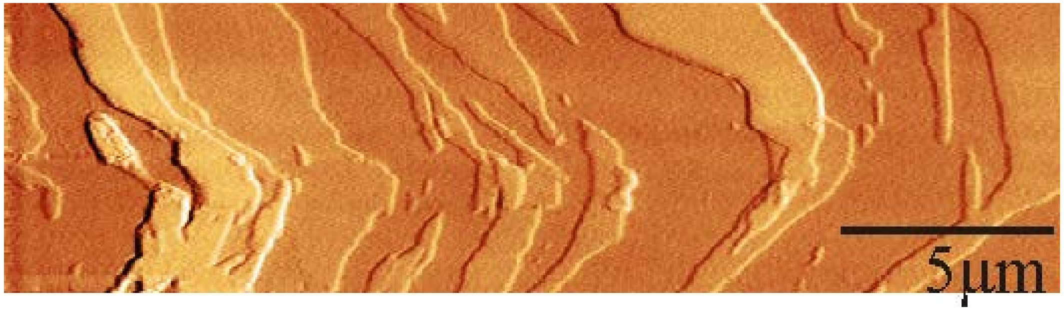
Content

Bi_2Te_3



Fabrication





Mechanical cleavage

Photolithography contacts

Hall measurements

$n=8.3 \times 10^{19}$ cm^{-3}

$\mu=250$ cm^2/Vs

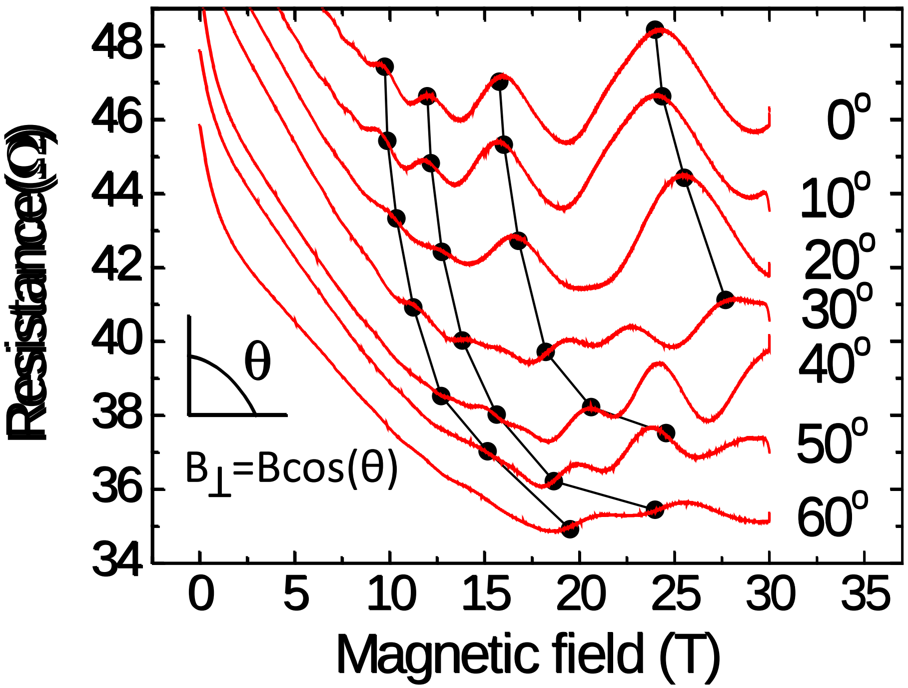
(bulk is conductive)

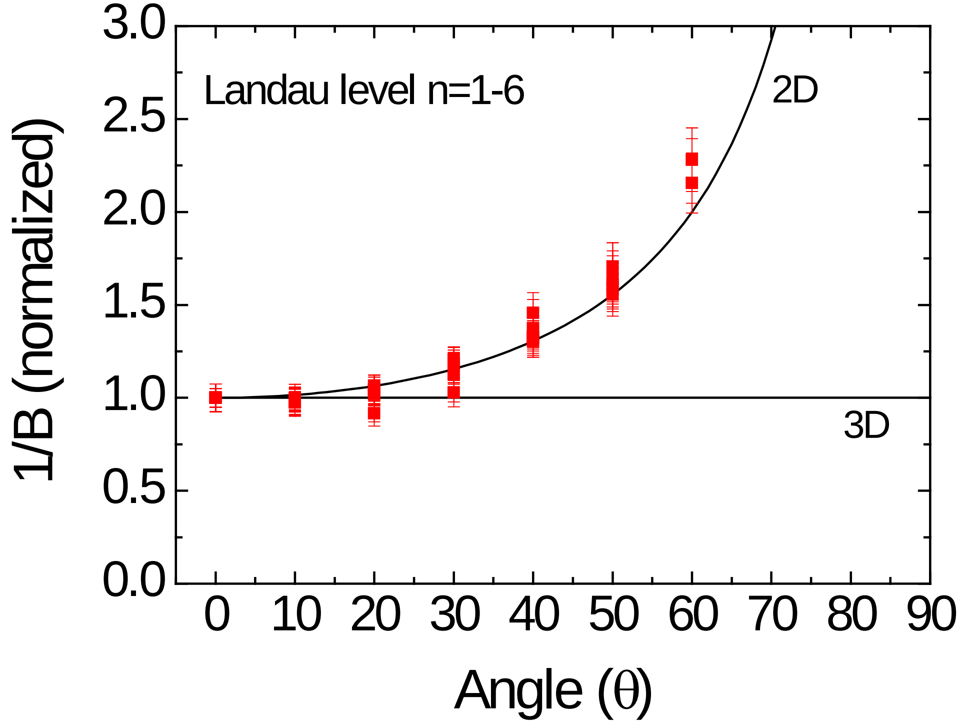
$l_{mfp}=22$ nm

~~$\mu B>>1$~~

So no quantum oscillations expected

Shubnikov-de-Haas oscillations

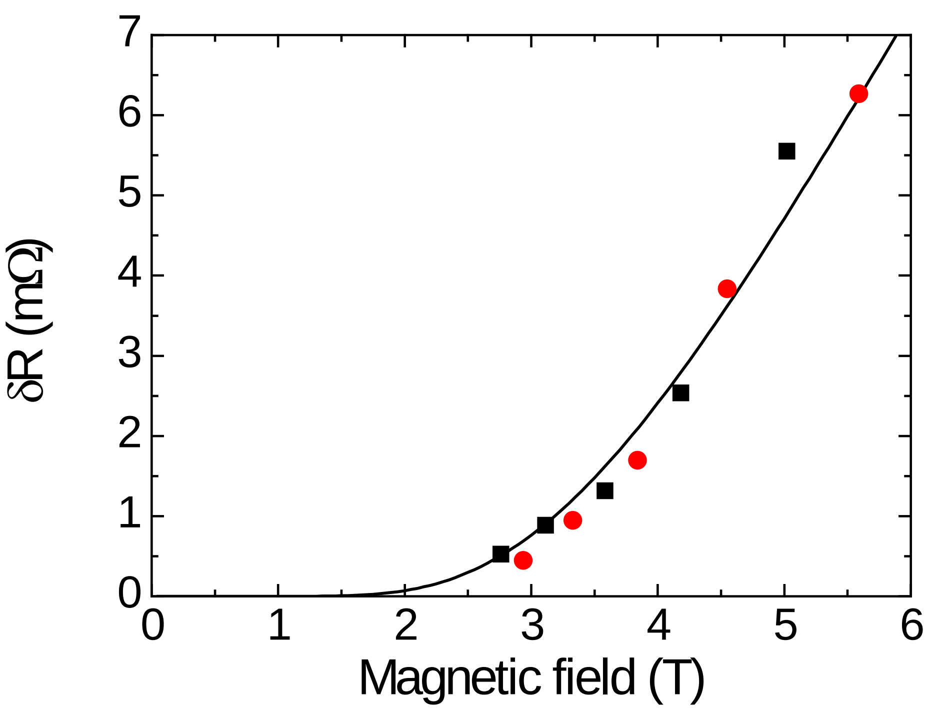


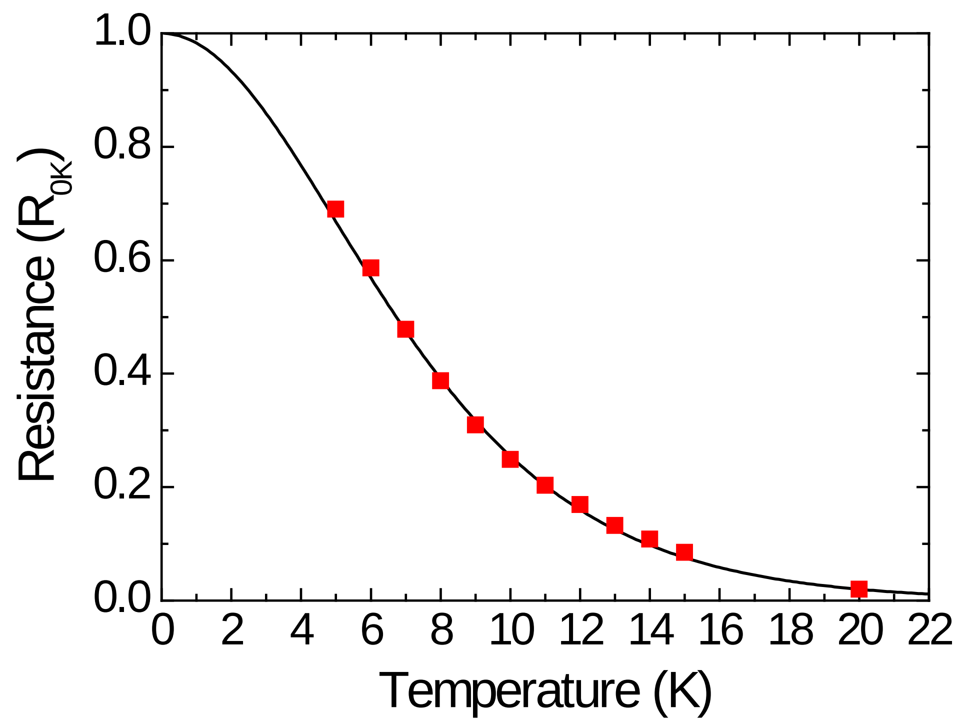


Left graph: Quantum oscillations @T=4.2K

Right graph: Oscillations from 2D channel

Shubnikov-de-Haas oscillations

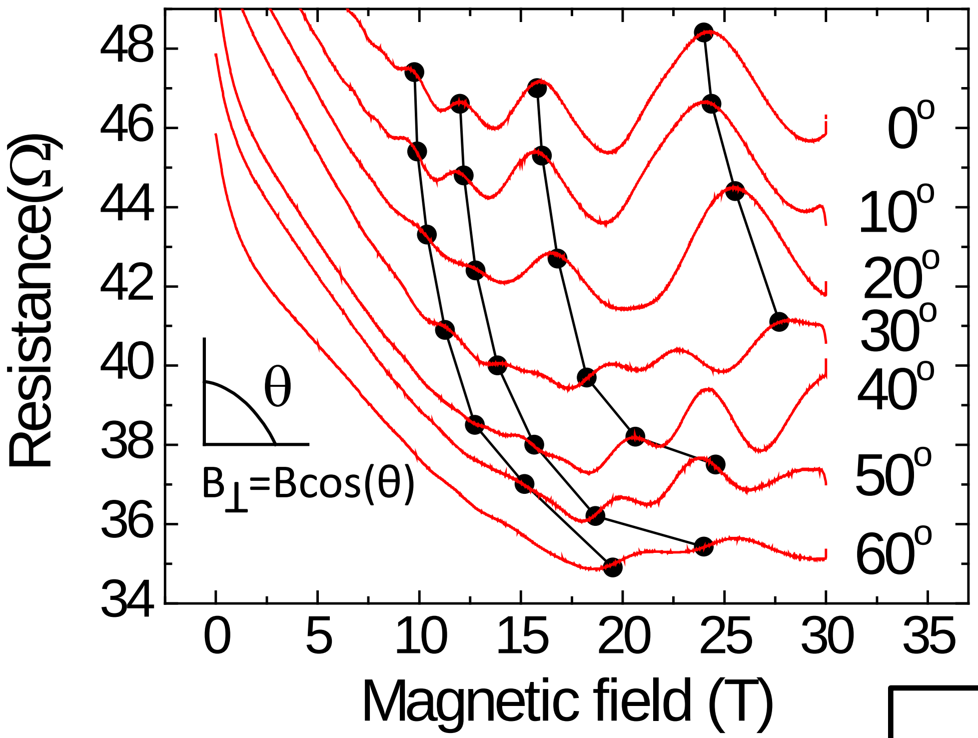




Left graph: Dingle temperature 1.65 K, μ=8300 cm^2/Vs

Right graph: Effective mass $0.16m_0$

Shubnikov-de-Haas oscillations



$$\delta\rho_{xx} \sim \cos\left(\frac{2\pi E_f}{\hbar\omega_c} + \pi + \varphi_B\right)$$

I.M. Lifshitz and A.M. Kosevich, Sov. Phys. JETP 2, **636** (1956)

Lifshitz-Kosevich formalism

Conductivity is the response function

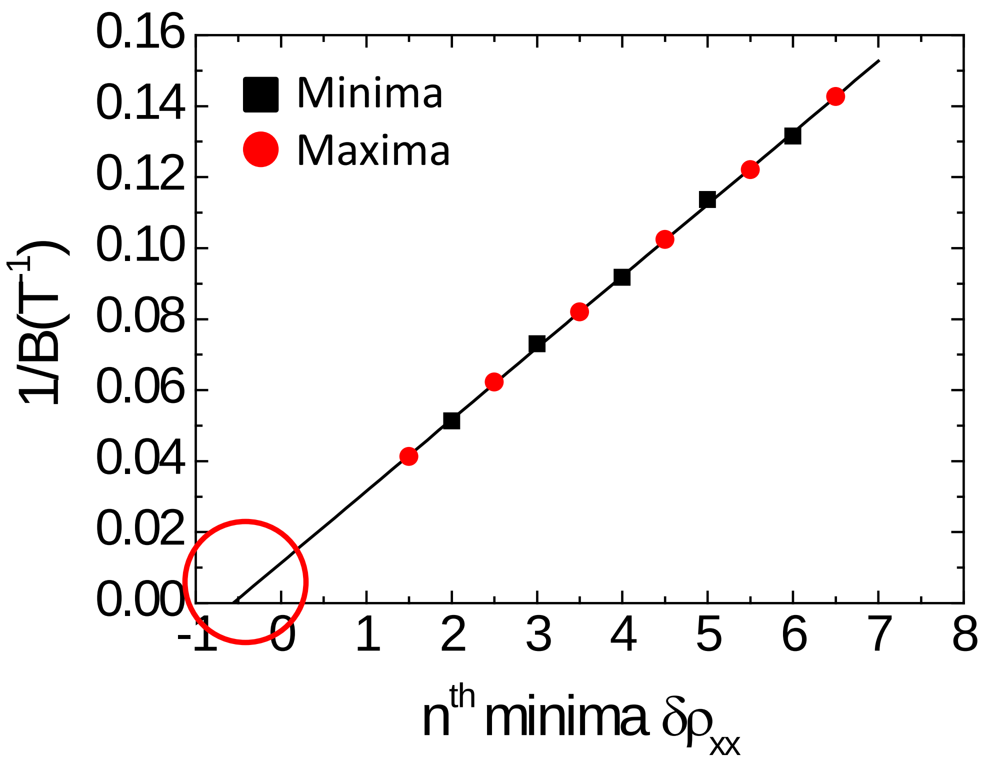
$$\rho_{xx} = \frac{\sigma_{xx}}{{\sigma_{xx}}^2 + {\sigma_{xy}}^2}$$

I.M. Lifshitz and A.M. Kosevich formalism applies if:

$\frac{\delta\sigma_{xx}}{\langle\sigma_{xx}\rangle} \ll 1$ or $\frac{\sigma_{xy}}{\sigma_{xx}} \gg 1$;Former is 0.01, later is 10

Then $\delta\rho_{xx} \sim \delta\sigma_{xx}$ So in normal case n^{th} minima=0 through 1/B=0

Lifshitz-Kosevich formalism

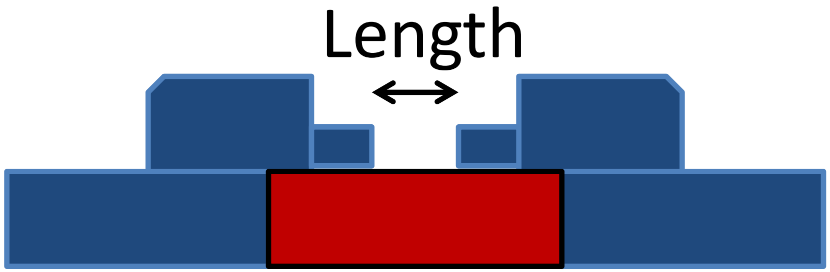


- Surface states present
- @ 1/B=0, n=-1/2 (Berry phase of π) → linear dispersion relation
- l_{mfp}=105 nm

Non-trivial surface states present

Content

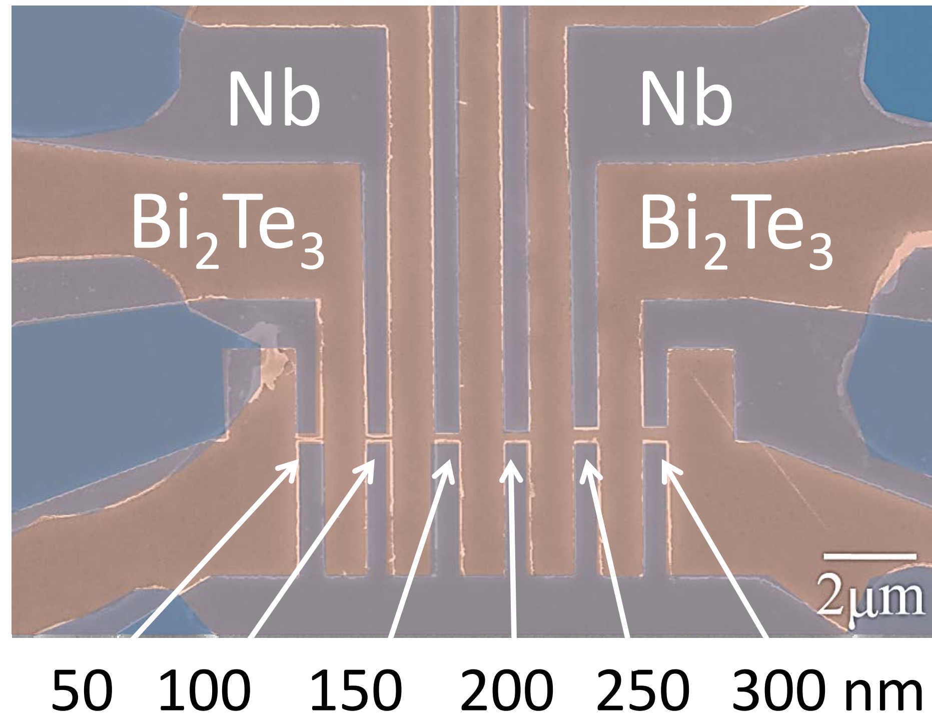
S/TI/S junctions



Nb: blue

Bi_2Te_3: red

Width=500 nm



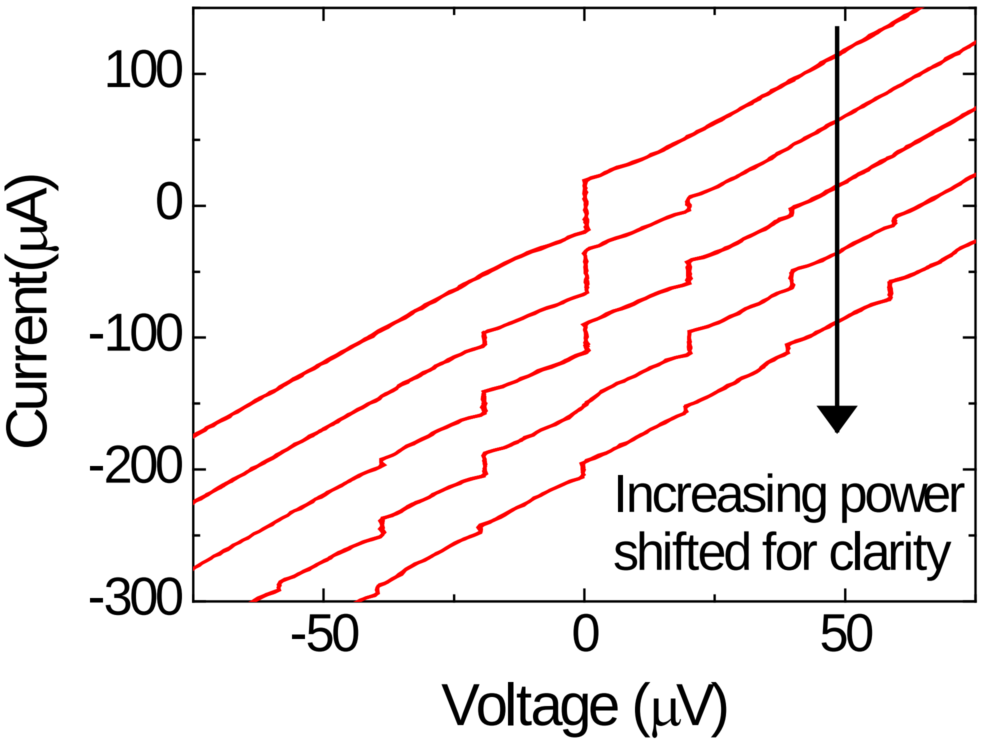
Junction lengths

Josephson supercurrent

Hallmarks for a Josephson junction:

1) Shapiro steps
2) Modulation I_c versus B-field

First hallmark – Shapiro steps



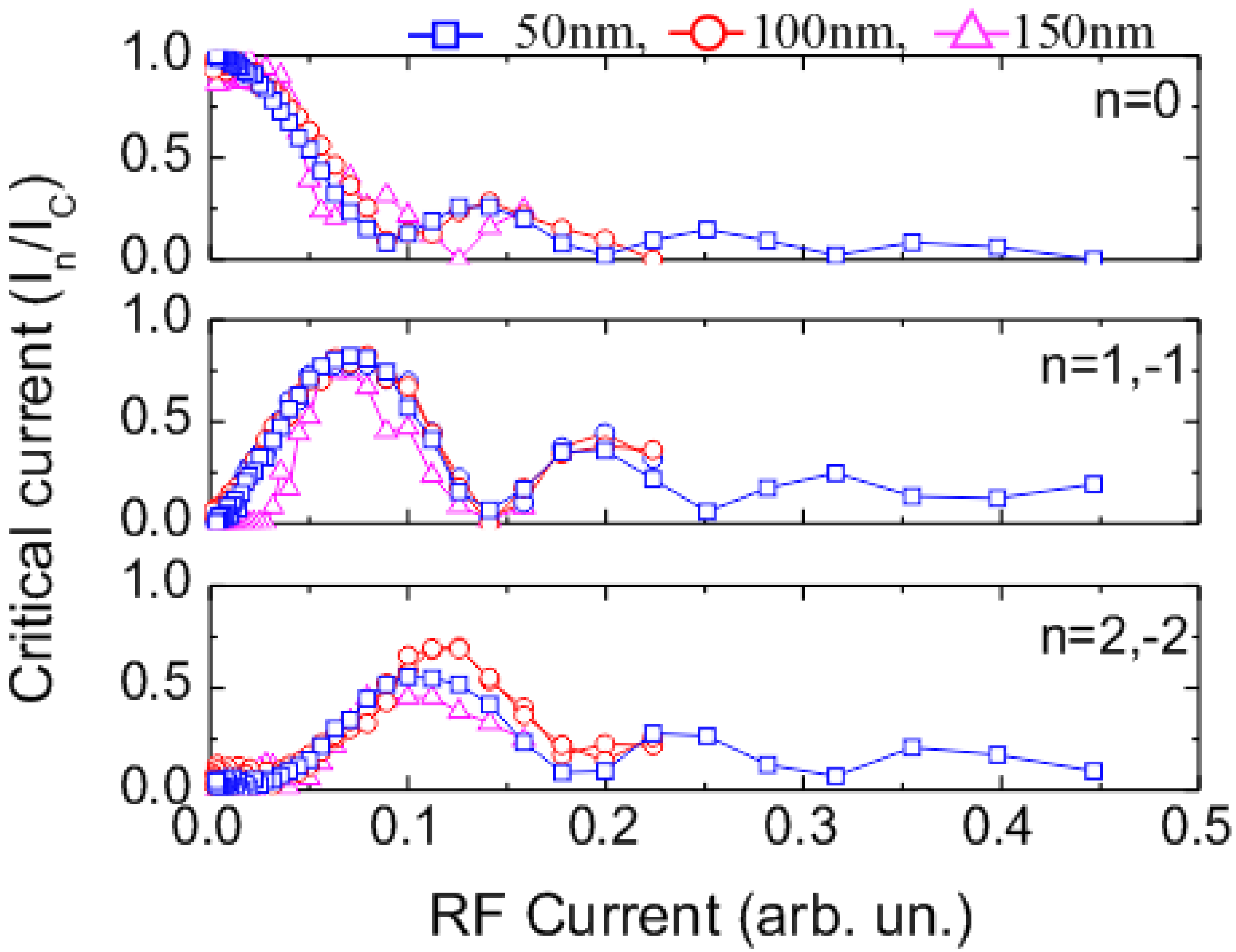
- Microwave frequency ω
- @$2eV/\hbar = n\omega$ Shapiro steps
- (Energy Cooper pairs resonant to energy microwave)

ω = 10 GHz

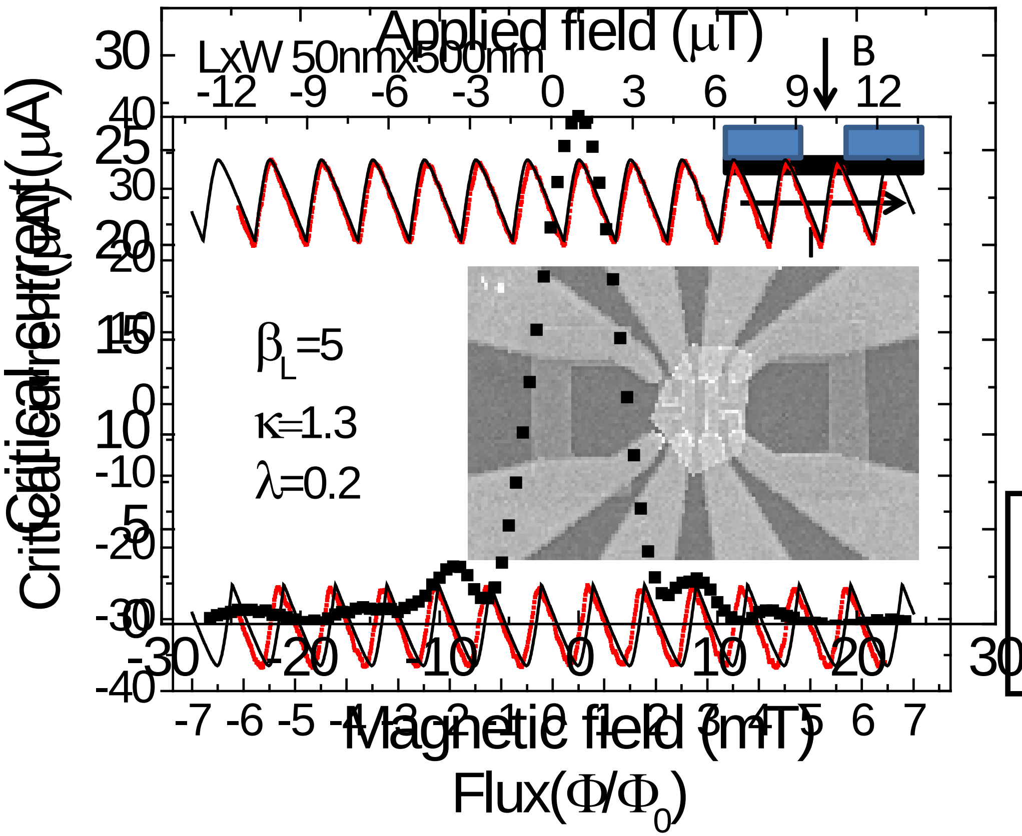
V= n x 20.7 μV

T=1.6 K

First hallmark – Shapiro steps



Second hallmark – I_c-B modulation

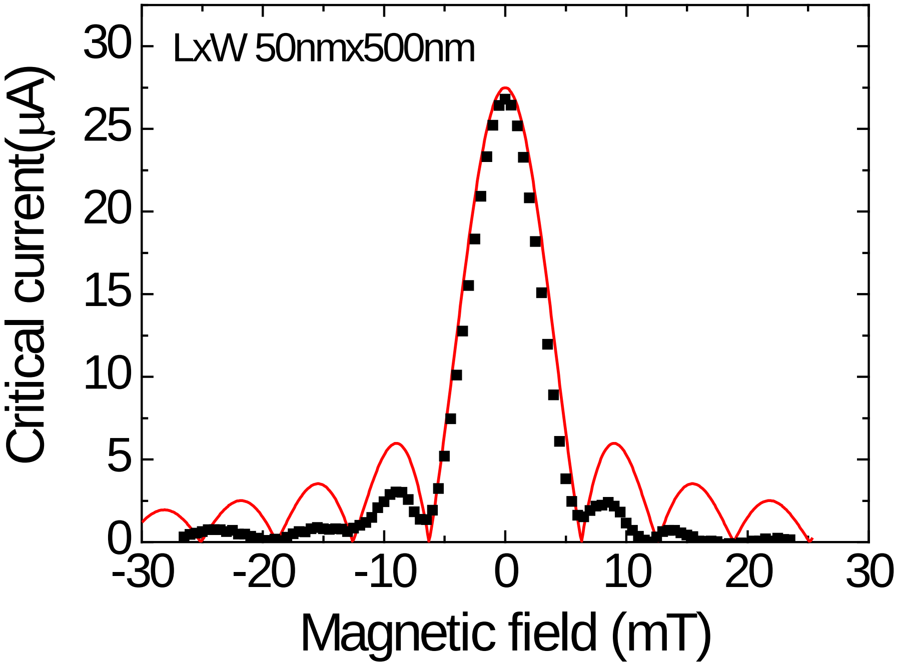


Modulation I_c

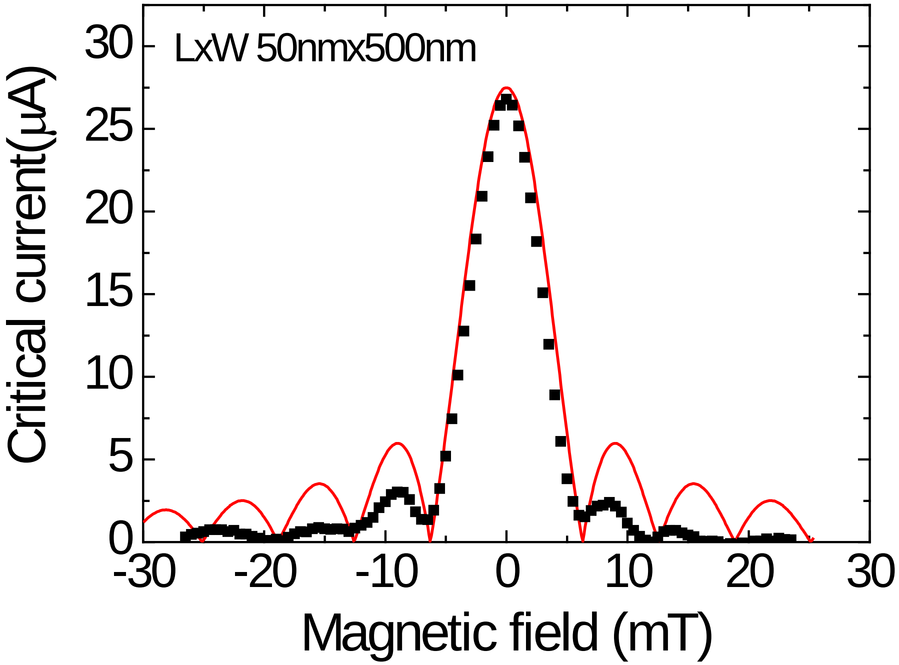
DC SQUID oscillations

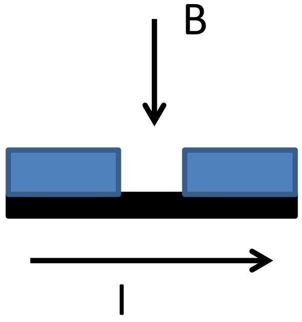
Josephson supercurrent present

Second hallmark – I_c-B modulation



Second hallmark – I_c-B modulation





Area uncertain:

- Penetration depth
- Flux focussing

Sinc function only valid for large L and W ratio

Content

Link supercurrent and surface states

Surface states

Supercurrent

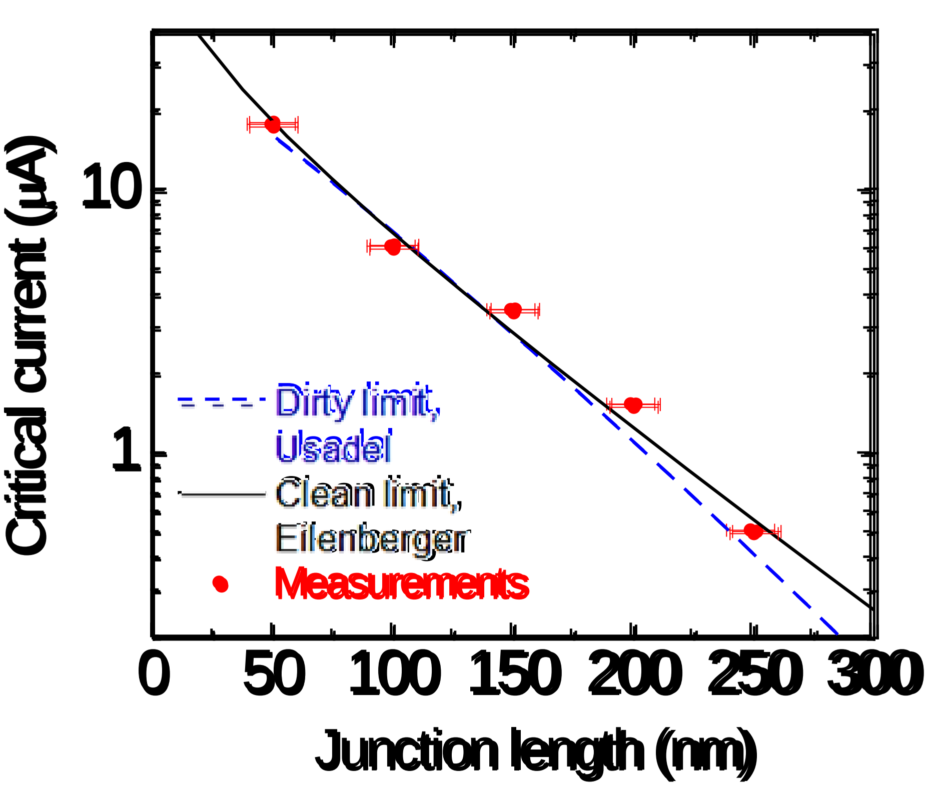
} Do we have a supercurrent through the surface states?

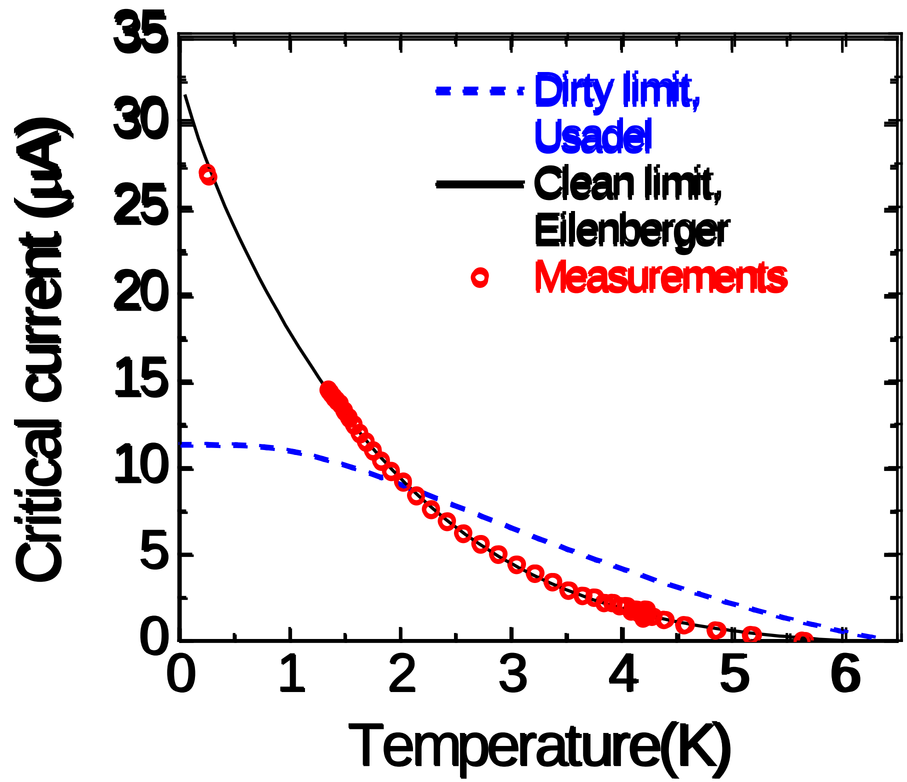
We have junctions between 50 and 250 nm and:

l_{mfp}=22 nm (bulk states) → diffusive transport

l_{mfp}=105 nm (surface states) → ballistic transport

Link supercurrent and surface states





Dirty or clean?

Definitely clean

Josephson supercurrent has been realized through the surface states of Bi_2Te_3

Provides prospects for Majorana devices……but

What is the best topological insulator for this purpose? (stability, insulating in the bulk, Dirac cone in the gap)

What is the smoking gun experiment with 3D topological insulators to observe Majorana fermions?

M. Veldhorst, M. Snelder et al. Nature Materials, 11, 417–421 (May 2012)